THE CLASSIFICATION OF SPINORS UNDER GSpin₁₄ OVER FINITE FIELDS

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ABSTRACT. The spinors of a 14-dimensional vector space V are studied with respect to the group $\operatorname{GSpin}_{14}$ of the 14-dimensional vector space V over finite fields \mathbf{F}_q . Results are given as follows: (1) the decomposition of the space of spinors into $\operatorname{GSpin}_{14}$ -equivalence classes or "orbits" over \mathbf{F}_q , (2) the structure of the fixer of $\operatorname{GSpin}_{14}$ for each orbit as an \mathbf{F}_q -group.

Introduction

Between 1969 and 1975, Atiyah, Bernstein, and S. I. Gel'fand [1, 2] proved the following theorem: If K is a local field of characteristic zero and $f(x) \in K[x_1, \ldots, x_n] - \{0\}$, then $|f|^s$ has a meromorphic continuation to the whole s-plane, where $|f|^s$ is a distribution in K^n , called the "complex power" of f(x) defined as

$$|f|^s(\mathbf{\Phi}) = \int_{K^n} |f(x)|_K^s \mathbf{\Phi}(x) \, dx,$$

where $|\cdot|_K$ is an absolute value in K, Φ is a Schwartz-Bruhat function, dx is a Haar measure on K^n , and s is a complex parameter restricted to the right-half plane. Furthermore, the candidates for the poles of $|f|^s$ can be written in terms of the roots of the b-function, or the Bernstein-Sato polynomial b(s) of f(x), i.e., if

$$b(s) = \prod_{\lambda > 0} (s + \lambda),$$

then the candidates for the poles of $|f|^s$ are $-\lambda$, $-\lambda-1$, $-\lambda-2$, It is also known that the λ 's in $b(s) = \prod_{\lambda>0} (s+\lambda)$ are positive rational numbers (cf. Kashiwara [14]). When K is a p-adic local field with q as the cardinality of its residue field, Igusa [6] has shown that $|f|^s$ is a rational function of $t=q^{-s}$. Furthermore, many examples suggest that an intimate relation between the real poles of $|f|^s$ and the roots of b(s), similar to the one in the archimedean case, would also exist in the p-adic case. Recently, Loeser [17] proved that the real poles of $|f|^s$ are the roots of b(s) if n=2.

Partly because the p-adic case of the theory of the complex powers is not as satisfactory as the archimedean case, Igusa began to study a certain complex-valued p-adic integral especially for the "regular prehomogeneous vector space"

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cases (cf. Sato, Kimura, and Shintani [15, 22]. Serre [23] named it "Igusa's local zeta funcion."

A regular prehomogeneous vector space defined over a field F is a triplet (G, ρ, X) in which $X = \operatorname{Aff}^n$ for some n, G is a reductive F-subgroup of $\operatorname{GL}(X) = \operatorname{GL}_n$ acting transitively on the complement Y of an absolutely irreducible F-hypersurface f(x) = 0 in X, and ρ is a rational representation of G on X. f(x) is a homogeneous relative G-invariant, i.e., $f(gx) = \nu(g)f(x)$ for every g in G, and is unique up to a factor in $F - \{0\}$. If K is a p-adic completion of a number field, O_K is the ring of integers of K, $\pi_K O_K$ is the maximal ideal of O_K , and q is the cardinality of $O_K/\pi_K O_K$, then the Igusa local zeta function is defined for any f(x) in $K[x_1, \ldots, x_n] - \{0\}$ and for $\operatorname{Re}(s) > 0$ as

$$Z(s) = \int_{O_{\nu}^{n}} |f(x)|_{K}^{s} dx = |f|^{s} (\Phi_{0}),$$

where Φ_0 is the characteristic function of $X(O_K)$ and dx is the Haar measure on K^n normalized as $\operatorname{vol}(O_K^n) = 1$. Igusa [6] has shown that Z(s) has a meromorphic continuation to the whole s-plane and is a rational function of $t = q^{-s}$. Z(s) has been computed for twenty of twenty-nine types of irreducible regular prehomogeneous vector spaces (cf. Igusa [12]).

Due to the difficulty in determining the local zeta function for some of the group invariants, Igusa [5, 7-11] developed a series of methods to calculate the local zeta function as the summation of the product of the cardinalities of the group orbits and the corresponding local zeta functions, i.e.,

$$Z(s) = \int_{O_K^n} |f(x)|_K^s dx = \sum_{\xi \in R} |G(\mathbb{F}_q) \cdot \bar{\xi}| \int_{\xi + \pi O_K^n} |f(x)|_K^s dx.$$

Here R is a subset of O_K^n such that its image \overline{R} in \mathbb{F}_q^n forms a complete set of representatives of $G(\mathbb{F}_q)$ -orbits in \mathbb{F}_q^n , i.e.,

$$\coprod_{\xi \in R} G(\mathbf{F}_q) \cdot \bar{\xi} = \mathbf{F}_q^n.$$

It is therefore essential to understand the orbital structure of the group G over finite fields. $GSpin_{14}$ is one of the few algebaic groups associated with regular prehomogeneous vector spaces for which the Igusa local zeta function remains unknown. Based on the works of Popov [20] and Kac and Vinberg [13] concerning $Spin_{14}$ over an algebraically closed field of characteristic zero, the following results are obtained in this paper:

- (1) the decomposition of the space of spinors into $GSpin_{14}$ -equivalence classes or "orbits" over \mathbf{F}_q , and
- (2) the structure of the fixer of $GSpin_{14}$ for each orbit as an F_a -group.

One should point out that the conjecture on the relation between the real poles of the Igusa local function and the roots of the Bernstein-Sato polynimial is verified for $GSpin_{14}$ and the candidates for the real poles of the Igusa loca zeta function of $GSpin_{14}$ are determined in [16] by using the above classification. Futhermore, this conjecture has been verified for any reduced irreducible regular prehomogeneous vector spaces by T. Kimura, F. Sato, and X.-W. Zhu [16].

We should also mention that the poles of Z(s) for curves have been closely examined by D. Meuser [18, 19].

1. Preliminaries

If V is a finite 2n-dimensional vector space over a field k and if Q(u) is a nondegenerate quadratic form on V, we denote by $(\ ,\)$ the bilinear form associated to Q so that (u,v)=Q(u+v)-Q(u)-Q(v). We denote by C the Clifford algebra of the pair (V,Q) and by $x\to x'$ its canonical antiautomorphism. Then $C=C^+\oplus C^-$, where C^+ is the space of invariant elements with respect to this antiautomorphism and C^- is the space of anti-invariant elements.

The Clifford group is

$$\widetilde{G}^* = \{ s \in C ; s \text{ invertible in } C \text{ and } sVs^{-1} = V \}.$$

The even Clifford group is

$$(\widetilde{G}^*)^+ = \widetilde{G}^* \cap C^+.$$

The spin group is

$$\operatorname{Spin}_{2n} = \widetilde{G} = \{ s \in (\widetilde{G}^*)^+ ; ss' = 1 \}.$$

We denote by C_W the subalgebra of C generated by any subspace W of V. Note that if W is totally isotropic, i.e., if Q=0 on W, C_W is isomorphic to the exterior algebra $\Lambda(W)$ of W.

We shall use ϕ to denote the vector representation of Spin_{2n} which is the restriction of the epimorphism $\phi\colon \widetilde{G}^* \to \mathrm{Aut}(V\,,\,Q)$, given by $\phi(s)\cdot v = svs^{-1}$, to Spin_{2n} . This restriction is an epimorphism onto the connected component of identity of $\mathrm{Aut}(V\,,\,Q)$ with kernel $\{\pm 1\}$.

Let V = L + M, where L, M are maximal totally isotropic subspaces of V. Choose bases e_1, \ldots, e_n and e_{n+1}, \ldots, e_{2n} of L and M respectively satisfying $(e_i, e_{n+i}) = 1$ for $1 \le i \le n$ and $(e_i, e_j) = 0$ for any other pairs (i, j), $i \le j$. If $s \in \widetilde{G}^*$, we shall define four $n \times n$ matrices $\alpha, \beta, \gamma, \delta$ by

$$\phi(s) \cdot (e_1 \cdots e_{2n}) = (e_1 \cdots e_{2n}) \begin{bmatrix} \alpha & \beta \\ \gamma & \delta \end{bmatrix}.$$

Then, $\alpha^t \beta$, $\gamma^t \delta$ are alternating and $\alpha^t \delta + \beta^t \gamma = 1_n$ or, equivalently, ${}^t \alpha \gamma$, ${}^t \beta \delta$ are alternating and ${}^t \alpha \delta + {}^t \gamma \beta = 1_n$. We shall write

$$\phi(s) = \begin{bmatrix} \alpha & \beta \\ \gamma & \delta \end{bmatrix}.$$

 C^+ is an enveloping algebra of \widetilde{G} and it is isomorphic to the direct sum of two total matrix algebras of degree 2^{n-1} . The representation ρ of \widetilde{G} in C^+ is called the spin representation of \widetilde{G} . It is the direct sum of two irreducible halfspin representations of \widetilde{G} of degree 2^{n-1} . An element of the representation space is called a spinor. We shall make the half-spin representation explicit as follows.

Put $e_L = e_1 \cdots e_n$ and $e_M = e_{n+1} \cdots e_{2n}$. Then Ce_M is a minimal left ideal of C and the correspondence $x \to xe_M$ gives an isomorphism $C_L \to Ce_M$ of

these vector spaces. Therefore there exists a unique element y of C_L satisfying $sxe_M = ye_M$ for any $s \in C$ and $x \in C_L$. Let

$$X = (C_L)^+ = C_L \cap C^+ \cong \Lambda^+(L) = \Lambda^+(Aff^n),$$

i.e., the sum of all even degree homogeneous parts of $\Lambda(L)$. This is a vector space of dimension 2^{n-1} defined over k. By restricting ρ to \widetilde{G} , we get a half-spin representation of \widetilde{G} in X.

Let $s_i(\lambda) = \lambda^{-1} + (\lambda - \lambda^{-1})e_i e_{n+i}$ with $\lambda \in G_m$ and $1 \le i \le n$. Then $\lambda \to s_i(\lambda)$ gives a homomorphism $G_m \to \widetilde{G}$. Since the one-parameter subgroups $\{s_i(\lambda)\}$ and $\{s_j(\lambda)\}$ commute, the mapping $(G_m)^n \to \widetilde{G}$ defined by $(t_1, \ldots, t_n) \to \prod_{i=1}^n s_i(t_i)$ is a homomorphism. Its image group \widetilde{T} is a maximal torus of \widetilde{G} and its kernel consists of $(\pm 1, \ldots, \pm 1)$ with an even number of minus signs.

Let $s_{ij}(\lambda) = 1 + \lambda e_i e_j$ with $\lambda \in G_a$ and $1 \le i \ne j \le 2n$ such that $(e_i, e_j) = 0$. Then the correspondence $\lambda \to s_{ij}(\lambda)$ defines an isomorphism of G_a to its image group P_{ij} in \widetilde{G} . These image groups are the one-parameter unipotent subgroups of \widetilde{G} corresponding to the 2n(n-1) roots of \widetilde{G} relative to \widetilde{T} .

If W is a subspace of V and if $\Lambda^2(W)$ denotes the homogeneous part of degree two in the exterior algebra C_W , we get a well-defined homomorphism $\exp: \Lambda^2(W) \to \widetilde{G}$ such that $\exp(x) = 1 + x + \frac{1}{2!}x^2 + \cdots$ for every x in $\Lambda^2(W)$. The element $\exp(x)$ is called the exponential of x.

Let

$$G = \operatorname{GSpin}_{2n} = (\widetilde{G}^*)^+.$$

Then

$$g_i(t) = t^{1/2} s_i(t)^{1/2} = 1 + (t-1)e_i e_{n+i}, \qquad 1 \le i \le n$$

are elements of $GSpin_{2n}$ for any $t \in G_m$ and

$$T = \left\{ t = \prod_{i=1}^{n} g_i(t_i); \ t_1, \ldots, t_n \in G_m \right\}$$

is a maximal torus of G. We give the following easily proved lemma for later use.

Lemma. Let $1 \le i, j \le 7$. Then

$$\phi(s_{i,j}(\lambda)) \cdot e_k = e_k, \qquad k \neq n+i, n+j;$$

$$\phi(s_{i,j}(\lambda)) \cdot e_{n+i} = e_{n+i} - \lambda e_j;$$

$$\phi(s_{i,j}(\lambda)) \cdot e_{n+j} = e_{n+j} + \lambda e_i;$$

$$\phi(s_{n+i,n+j}(\lambda)) \cdot e_k = e_k, \qquad k \neq i, j;$$

$$\phi(s_{n+i,n+j}(\lambda)) \cdot e_i = e_i - \lambda e_{n+j};$$

$$\phi(s_{n+i,n+j}(\lambda)) \cdot e_j = e_j + \lambda e_{n+i};$$

$$\phi(s_{i,n+j}(\lambda)) \cdot e_k = e_k, \qquad k \neq n+i, j;$$

$$\phi(s_{i,n+j}(\lambda)) \cdot e_{n+i} = e_{n+i} - \lambda e_{n+j};$$

$$\phi(s_{i,n+j}(\lambda)) \cdot e_j = e_j + \lambda e_i,$$

where $\lambda \in G_a$. For any $t_1, \ldots, t_7 \in G_m$ we have

$$\phi\left(\prod_{i=1}^{7} g_{i}(t_{i})\right) \cdot e_{k} = \begin{cases} t_{k}e_{k}, & \text{if } 1 \leq k \leq 7, \\ t_{k}^{-1}e_{k}, & \text{if } 8 \leq k \leq 14, \end{cases}$$

$$\rho\left(\prod_{i=1}^{7} g_{i}(t_{i})\right) \cdot (e_{i_{1}} \cdots e_{i_{p}}) = (t_{i_{1}} \cdots t_{i_{p}})e_{i_{1}} \cdots e_{i_{p}}.$$

We shall use G_x to denote the fixer of G at the point $x \in X$ with respect to ρ , G_x^0 to denote the connected component of the identity of G_x , and Z(H) to denote the center of the group H. We define

 $e_{i_1\cdots i_p}^*$ = the partial products of e_1 , ..., e_n satisfying $e_{i_1}\cdots e_{i_p}e_{i_1\cdots i_p}^*=e_L$. In particular,

$$e_i^* = (-1)^{i-1}e_1 \cdots \hat{e}_i \cdots e_n,$$

$$e_{ij}^* = (-1)^{i+j-1}e_1 \cdots \hat{e}_i \cdots \hat{e}_j \cdots e_n,$$

$$e_{ijk}^* = (-1)^{i+j+k}e_1 \cdots \hat{e}_i \cdots \hat{e}_j \cdots \hat{e}_k \cdots e_n.$$

In the sequel we shall assume n = 7, $X = \Lambda^{+}(Aff^{7})$, and $G = GSpin_{14}$ unless the contrary is expressly stated.

2. Orbital decomposition

In this section we obtain the representatives of $GSpin_{14}$ -orbits in $\Lambda^+(Aff^7)$ over any finite field k of characteristic different from 2. Since most of the methods used in this section are the same as Popov's [20] for obtaining the representatives of $GSpin_{14}$ -orbits in $\Lambda^+(Aff^7)$ over algebraically closed fields of characteristic zero, we shall only sketch the procedure and examine those cases that lead to different results.

We shall say that elements in X are G-equivalent if they lie in the same orbit of G in X, i.e., $\rho(s) \cdot x = y$ for some $s \in G$, denoted by $x \xrightarrow{s} y$. For every element $u = \sum_{1 \le i < j \le 6} \alpha_{ij} e_{ij}^*$, $\alpha_{ij} \in k$, in $\Lambda^4(\operatorname{Aff}^6)$, the rank of the alternating matrix (α_{ij}) is called the rank of the spinor u and denoted as rank (u).

For any nonzero spinor $x \in X$, $x = x_0 + x_2 + x_4 + x_6$, where $x_0 \in k^*$, $x_i \in \Lambda^i(Aff^7)$ for i = 2, 4, 6. By Lemma 1 of [4], have

$$x \xrightarrow{s_1(x_0)} 1 + x_2' + x_4' + x_6' \xrightarrow{\exp(-x_2')} 1 + x_4'' + x_6''$$

Therefore, we can assume that $x = 1 + x_4 + x_6$. We shall study spinors of this type according to $x_6 \neq 0$ and $x_6 = 0$.

From Lemma 1 of [4], have

$$G_1 \cong \phi(G_1) = \left\{ \begin{bmatrix} \alpha & 0 \\ \gamma & \delta \end{bmatrix}, \delta = {}^t\!\alpha^{-1} \in SL_7, \ {}^t\!\alpha\gamma \in Alt_7 \right\}.$$

Let H be the subgroup of G_1 such that

$$\phi(H) = \left\{ \begin{bmatrix} \alpha & 0 \\ 0 & {}^t\alpha^{-1} \end{bmatrix}, \alpha \in SL_7 \right\},\,$$

let S be the subgroup of H for which

$$\alpha = \begin{bmatrix} \varepsilon & 0 \\ 0 & \det \varepsilon^{-1} \end{bmatrix}$$

with $\varepsilon \in GL_6$ in $\phi(S)$, and let \widetilde{S} be the subgroup of S for which $\varepsilon \in SL_6$ in $\phi(\widetilde{S})$. Thus $H \cong SL_7$, $S \cong GL_6$, and $\widetilde{S} \cong SL_6$.

First, consider the case when $x_6 \neq 0$. Then

$$x = 1 + x_4 + x_6 \xrightarrow{h \in H} 1 + x_4 + e_7^*$$

If $x_4 = 0$, then

$$(1) x = 1 + e_7^*.$$

If $x_4 \neq 0$, write x_4 in the form $x_4 = ye_7 + z$, where $y \in \Lambda^3(Aff^6)$, $z \in \Lambda^4(Aff^6)$. With respect to the group S, it is possible to bring x into one of the forms

- $(2) 1 + ye_7 + e_7^*,$
- $(3) 1 + ye_7 + e_{147}^* + e_7^*,$
- (4) $1 + ve_7 + e_{147}^* + e_{257}^* + e_7^*$, or

(5)
$$1 + ye_7 + e_{147}^* + e_{257}^* + e_{367}^* + pe_7^*, \qquad p \in k^*,$$

depending on whether rank(z) = 0, 2, 4, or 6, where $y \in \Lambda^3(Aff^6)$.

As in [20], we may claim that spinors of type (3) and (4) are equivalent to spinors of type (2). Futhermore, considering the action of S on spinors of type (2) and the fact that all the GL_6 -inequivalent trivectors of six-dimensional space are those GL_7 -inequivalent trivectors of seven-dimensional space involving six vectors, and by the classification of trivectors of seven-dimensional space over finite fields [11], it follows that y can be transformed to one of the pairwise S-inequivalent forms

$$\begin{split} e_1e_2e_3\,,\quad e_1e_2e_3+e_4e_5e_6\,,\quad e_1e_2e_3+e_3e_4e_5\,,\\ e_1e_2e_3+e_4e_3e_5+e_6e_5e_2\,,\quad \text{or}\\ 4^{-1}e_{123}^*+(4\lambda)^{-1}(e_{156}^*-e_{246}^*+e_{345}^*)\,,\qquad \lambda\in \mathbf{F}_q^*-(\mathbf{F}_q^*)^2. \end{split}$$

Therefore, with the action of S, x can be transformed into one of the following:

$$\begin{aligned} 1 + \nu e_1 e_2 e_3 e_7 + \nu^{-1} e_7^* \,, \\ 1 + \nu e_1 e_2 e_3 e_7 + \nu e_4 e_5 e_6 e_7 + \nu^{-1} e_7^* \,, \\ 1 + \nu (4^{-1} e_{123}^* + (4\lambda)^{-1} (e_{156}^* - e_{246}^* + e_{345}^*)) + \nu^{-1} e_7^* \,, \\ 1 + \nu e_1 e_2 e_3 e_7 + \nu e_3 e_4 e_5 e_7 + \nu^{-1} e_7^* \,, \text{ or } \\ 1 + \nu e_1 e_2 e_3 e_7 + \nu e_4 e_3 e_5 e_7 + \nu e_6 e_5 e_2 e_7 + \nu^{-1} e_7^* \,, \end{aligned}$$

where $\nu \in k^*$. Applying the appropriate $h \in H$ such that $\phi(h)e_i = \lambda_i e_i$ with $\lambda_i \in k^*$, $1 \le i \le 6$, to the above spinors, we conclude that spinors of type (2) are equivalent to one of the following spinors:

- (6) $1 + e_1 e_2 e_3 e_7 + e_7^*$,
- (7) $p(1+e_1e_2e_3e_7+e_4e_5e_6e_7+e_7^*), p \in k^*,$
- (8) $p(1+4^{-1}e_{123}^*+(4\lambda)^{-1}(e_{156}^*-e_{246}^*+e_{345}^*)-e_7^*), p \in k^*,$
- (9) $1 + e_1e_2e_3e_7 + e_3e_4e_5e_7 + e_7^*$, or
- $(10) 1 + e_1e_2e_3e_7 + e_4e_3e_5e_7 + e_6e_5e_2e_7 + e_7^*,$

where $\lambda \in \mathbf{F}_q^* - (\mathbf{F}_q^*)^2$, with q a power of any odd number.

Now we consider the spinor of type (5). Let R be the subgroup of \widetilde{S} such that

$$\phi(R) = \left\{ \begin{bmatrix} A & 0 \\ 0 & {}^tA^{-1} \end{bmatrix} \; ; \; A = \begin{bmatrix} \varepsilon & 0 \\ 0 & 1 \end{bmatrix} \; , \; \varepsilon \in SP_6 \right\}.$$

Then $R \cong SP_6$. By Lemma 5 of [4], $R \subset G_{a_0}$ for $a_0 = e_1e_4 + e_2e_5 + e_3e_6$. It is also clear that $\Lambda^p(Aff^7)$, $\Lambda^p(Aff^6)$, and spinor e_7 remain invariant under R. Hence, spinors of type (5), $x = 1 + ye_7 + e_{147}^* + e_{257}^* + e_{367}^* + pe_7^*$, are R-inequivalent if and only if $y \in \Lambda^3(Aff^6)$ are SP_6 -inequivalent. We shall first consider the classification of the trivector of six-dimensional space with respect to SP_6 .

Theorem 1. Every trivector in $\Lambda^3(Aff^6)$ is SP_6 -equivalent to one of the following twenty pairwise SP_6 -inequivalent trivectors:

- (a) 0;
- (b) $e_1e_2e_3$;
- (c) $e_1e_2e_5 + e_1e_3e_6$;
- (d) $e_1e_2e_3 + pe_4e_5e_6$, $p \in k^*$;
- (e) $e_1e_4e_3 + e_5e_2e_3 + p(e_3e_1e_4 + e_3e_2e_5), p \in k^*$;
- (f) $e_1e_2e_3 + pe_4e_5e_6 + e_4e_2e_5 + e_4e_3e_6$, p = 0, $p \in k^*$;
- (g) $e_1e_4e_3 + e_5e_2e_3 + e_1e_2e_6 + e_4e_2e_5 + e_4e_3e_6$;
- (h) $e_1e_2e_3 + pe_4e_5e_6 + e_1e_2e_5 + e_1e_3e_6 + e_5e_1e_4 + e_5e_3e_6$, with $p \in k^*$;
- (i) $e_1e_4e_3 + e_5e_2e_3 + e_1e_2e_5 + e_1e_3e_6 + p(e_2e_1e_4 + e_2e_3e_6)$, with p = 0, p = 1;
- (j) $e_1e_4e_3 + e_5e_2e_3 + pe_1e_2e_6 + r(e_6e_1e_4 + e_6e_2e_5)$, with

$$\begin{cases} p=0 \\ r=0 \end{cases}, \quad \begin{cases} p=0 \\ r \in k^* \end{cases}, \quad \begin{cases} p=1 \\ r=0 \end{cases}, \quad \begin{cases} p=1 \\ r \in k^* \end{cases}$$

(k) $e_1e_2e_3 + re_4e_5e_6 + e_1e_2e_5 + e_1e_3e_6 + p(e_4e_2e_5 + e_4e_3e_6)$, with

$$\begin{cases} p=0 \\ r=0 \end{cases}, \; \begin{cases} p=0 \\ r\in k^* \end{cases}, \; \begin{cases} p\in k^* \\ r\in k^* \end{cases};$$

- (1) $2^{-1}e_2e_3e_4 + (2\lambda)^{-1}e_1e_3e_5$;
- (m) $2^{-1}e_2e_3e_4 + (2\lambda)^{-1}e_1e_3e_5 + 2e_1e_2e_5 + 2e_1e_3e_6$,

where $\lambda \in \mathbf{F}_q^* - (\mathbf{F}_q^*)^2$, with q a power of any odd prime number.

Proof. The first eighteen pairwise SP_6 -inequivalent trivectors (a)-(k) are obtained in the same way as in [20]. However the fixers of $\xi = e_1e_4e_3 + e_5e_2e_3$ and $\eta = e_1e_4e_3 + e_5e_2e_3 + e_1e_2e_5 + e_1e_3e_6 + e_2e_1e_4 + e_2e_3e_6$ in SP_6 have two connected components, and the representatives ξ' and η' of the other connected components are the elements of SP_6 defined by $e_1 \rightarrow e_2$, $e_2 \rightarrow e_1$, $e_3 \rightarrow -e_3$, i.e.,

$$\phi(\xi')\,,\,\phi(\eta') = \begin{bmatrix} \alpha & 0 \\ 0 & {}^t\alpha^{-1} \end{bmatrix}\,, \qquad \alpha = \begin{bmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix}.$$

Take λ from $\mathbf{F}_q^* - (\mathbf{F}_q^*)^2$, and put $\mu = \lambda^{1/2}$ so that $\mu^{\sigma} = -\mu$ for $\sigma \in \text{Gal}(\mathbf{F}_{q^2}/\mathbf{F}_q)$. Choose $g \in SP_6(\mathbf{F}_{q^2})$ such that

$$g = \begin{bmatrix} g_1 & 0 \\ 0 & {}^tg_1^{-1} \end{bmatrix}, \qquad g_1 = \begin{bmatrix} 1 & 1 & 0 \\ \mu & -\mu & 0 \\ 0 & 0 & -(2\mu)^{-1} \end{bmatrix} \in SL_3,$$

and $\gamma \in G_{\xi}$ such that

$$\phi(\gamma) = \begin{bmatrix} \alpha & 0 \\ 0 & {}^t \alpha^{-1} \end{bmatrix}, \qquad \alpha = \begin{bmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix}.$$

Then $G_{\xi}=G_{\xi}^0 \coprod \gamma G_{\xi}^0$, $\gamma^2=1$. By a theorem of Igusa's [11], we have $g^{\sigma}=g\gamma$ and

$$\begin{split} \xi' &= g \cdot (e_1 e_4 e_3 + e_5 e_2 e_3) \\ &= (e_1 + \mu e_2)(2^{-1} e_4 + (2\mu)^{-1} e_5)(-(2\mu)^{-1} e_3) \\ &+ (2^{-1} e_4 - (2\mu)^{-1} e_5)(e_1 - \mu e_2)(-(2\mu)^{-1} e_3) \\ &= 2^{-1} e_2 e_3 e_4 + (2\lambda)^{-1} e_1 e_3 e_5 \,, \\ \eta' &= g \cdot (e_1 e_4 e_3 + e_5 e_2 e_3 + e_1 e_2 e_5 + e_1 e_3 e_6 + e_2 e_1 e_4 + e_2 e_3 e_6) \\ &= 2^{-1} e_2 e_3 e_4 + (2\lambda)^{-1} e_1 e_3 e_5 + (e_1 + \mu e_2 + e_1 - \mu e_2)(-(2\mu)^{-1} e_3)(-2\mu e_6) \\ &+ (e_1 + \mu e_2)(e_1 - \mu e_2)(2^{-1} e_4 - (2\mu)^{-1} e_5 - 2^{-1} e_4 - (2\mu)^{-1} e_5) \\ &= 2^{-1} e_2 e_3 e_4 + (2\lambda)^{-1} e_1 e_3 e_5 + 2e_1 e_2 e_5 + 2e_1 e_3 e_6 \,, \end{split}$$

which are the trivectors (1) and (m).

Now if the factor y in spinor (5) has one of the forms in Theorem 1, other than type (a), (b), and (l), then it is equivalent to the spinors of the form

$$1 + ve_7 + z + pe_7^*$$

with $p \in k^*$, $y \in \Lambda^3(Aff^6)$, $z \in \Lambda^4(Aff^6)$, and rank(z) < 6 by the action of $h \in H$ such that $\phi(h)e_i = e_i$ for $1 \le i \le 6$, and

$$\phi(h)e_7 = \begin{cases} e_4 + e_7, & \text{if } y = \text{types (c), (i), (k);} \\ -e_1 + e_7, & \text{if } y = \text{types (f), (g);} \\ -e_2 + e_7, & \text{if } y = \text{type (h);} \\ e_6 + e_7, & \text{if } y = \text{type (d);} \\ -p^{-1}e_2 + e_5 + e_7, & \text{if } y = \text{type (e) and } p \neq -1; \\ -2^{-1}e_6 + e_7, & \text{if } y = \text{type (e) and } p = -1; \\ 2^{-1}e_4 + e_7, & \text{if } y = \text{type (m).} \end{cases}$$

Consequently, it is equivalent to one of the spinors of types (6)–(10). Otherwise we have the following situation:

(a) Spinors of type (5) with y = 0 are equivalent to the spinor

$$1 + e_{147}^* + e_{257}^*$$

when $p = \pm 2$, and

$$(12) 1 + (\sqrt{p^2 - 4})e_7^*$$

when $p \neq \pm 2$, over $k(\sqrt{p^2-4})$.

(b) Spinors of type (5) with $y = e_1 e_2 e_3$ are equivalent to the spinor

$$(13) 1 + e_1 e_2 e_3 e_7 + e_{147}^* + e_{257}^*$$

when $p = \pm 2$, and

$$(14) 1 + e_1 e_2 e_3 e_7 + (\sqrt{p^2 - 4}) e_7^*$$

when $p \neq \pm 2$, over $k(\sqrt{p^2-4})$.

(1) Spinors of type (5) with $y = 2^{-1}e_2e_3e_4 + (2\lambda)^{-1}e_1e_3e_5$, $\lambda \in \mathbf{F}_q^* - (\mathbf{F}_q^*)^2$, are equivalent to spinor (9) over $k(\sqrt{p^2-4})$ when $p \neq \pm 2$, and

$$(15) 1 + 4^{-1}e_{123}^* + (4\lambda)^{-1}(e_{156}^* - e_{246}^* + e_{345}^*)$$

when $p = \pm 2$.

Proof. (a), (b) Let $\varepsilon = p/2$. Then

 $1 + e_{147}^* + e_{257}^* + e_{367}^* + pe_7^*$

$$1 + e_{147}^* + e_{257}^* + e_{367}^* + pe_7^*$$

$$\xrightarrow{s_{1,4}(-\varepsilon)s_{2,5}(-\varepsilon)s_{10,13}(-\varepsilon)} 1 + e_{147}^* + e_{257}^* + \left(1 - \frac{p^2}{4}\right)e_{367}^*$$

$$= 1 + e_{147}^* + e_{257}^*, \quad \text{when } p = \pm 2,$$

and

$$1 + e_{1}e_{2}e_{3}e_{7} + e_{147}^{*} + e_{257}^{*} + e_{367}^{*} + pe_{7}^{*}$$

$$\xrightarrow{s_{1,4}(-\varepsilon)s_{2,5}(-\varepsilon)s_{10,13}(-\varepsilon)} 1 + e_{1}e_{2}e_{3}e_{7} + e_{147}^{*} + e_{257}^{*} + \left(1 - \frac{p^{2}}{4}\right)e_{367}^{*}$$

$$= 1 + e_{1}e_{2}e_{3}e_{7} + e_{147}^{*} + e_{257}^{*}, \quad \text{when } p = \pm 2.$$

When $p \neq \pm 2$, let $\omega = (p - \sqrt{p^2 - 4})/2$. Then

$$\frac{s_{9,12}(\omega)s_{8,11}(\omega)}{+(1-\omega p)(e_{147}^* + e_{257}^*) + e_{367}^* + pe_7^*} + (1-\omega p)(e_{147}^* + e_{257}^*) + e_{367}^* + pe_7^* \\
\frac{s_1(1-\omega^2)}{+(1-\omega^2)} + \omega(2-\omega p)(1-\omega^2)e_3e_6 + \omega(1-\omega^2)^{-1}e_2e_5 \\
+ \omega(1-\omega^2)e_1e_4 + (1-\omega p)(1-\omega^2)^{-1}e_{147}^* \\
+ (1-\omega p)(1-\omega^2)e_{257}^* + (1-\omega^2)e_{367}^* + p(1-\omega^2)e_7^* \\
\frac{\exp(-\omega(2-\omega p)(1-\omega^2)^{-1}e_3e_6)}{1-\omega^2} + \omega(1-\omega^2)^{-1}e_2e_5 + \omega(1-\omega^2)e_1e_4 \\
+ \left[\frac{1-\omega p}{1-\omega^2} + \frac{\omega^2(2-\omega p)}{(1-\omega^2)^2}\right]e_{147}^* \\
+ [(1-\omega p)(1-\omega^2) + \omega^2(2-\omega p)]e_{257}^* \\
+ (1-\omega^2)e_{367}^* + (p-2\omega)e_7^*$$

$$= 1 + \omega(1 - \omega^2)^{-1}e_2e_5 + \omega(1 - \omega^2)e_1e_4 + (1 - \omega^2)e_{367}^* + (p - 2\omega)e_7^*$$

$$\xrightarrow{\exp(-\omega(1-\omega^2)^{-1}e_2e_5 - \omega(1-\omega^2)e_1e_4)} 1 + \omega_1e_{367}^* + \omega_2e_7^*,$$

$$\omega_1 = 1 + \omega^2$$
, $\omega_2 = p - 2\omega = \sqrt{p^2 - 4}$

$$\xrightarrow{s_{10,13}(\omega_1\omega_2^{-1})} 1 + (\sqrt{p^2 - 4})e_7^*.$$

Applying the same procedure as above to $1 + e_1e_2e_3e_7 + e_{147}^* + e_{257}^* + e_{367}^* + pe_7^*$, followed by an application of the action of $h \in H$ such that $\phi(h)e_i = e_i$ for

 $i \neq 3, 6, \ \phi(h)e_3 = (1-\omega^2)^{-1}e_3$, and $\phi(h)e_6 = (1-\omega^2)e_6$ yields $1 + e_1e_2e_3e_7 + (\sqrt{p^2-4})e_7^*$.

(1) Let $\varepsilon = p/2$. Then

$$1 + 2^{-1}e_{2}e_{3}e_{4}e_{7} + (2\lambda)^{-1}e_{1}e_{3}e_{5}e_{7} + e_{147}^{*} + e_{257}^{*} + e_{367}^{*} + pe_{7}^{*}$$

$$\xrightarrow{s_{1,4}(-\varepsilon)s_{2,5}(-\varepsilon)s_{10,13}(-\varepsilon)} 1 + 2^{-1}e_{2}e_{3}e_{4}e_{7} + (2\lambda)^{-1}e_{1}e_{3}e_{5}e_{7}$$

$$+ e_{147}^{*} + e_{257}^{*} + \left(1 - \frac{p^{2}}{4}\right)e_{367}^{*}.$$

When $p = \pm 2$, consider $h \in H$ such that

$$\phi(h)e_1 = -e_1$$
, $\phi(h)e_2 = -e_4$, $\phi(h)e_3 = (8\lambda)^{-1}e_7$, $\phi(h)e_4 = e_2$, $\phi(h)e_5 = \lambda e_5$, $\phi(h)e_6 = -2e_6$, $\phi(h)e_7 = -4e_3$.

Then

$$1 + 2^{-1}e_{2}e_{3}e_{4}e_{7} + (2\lambda)^{-1}e_{1}e_{3}e_{5}e_{7} + e_{147}^{*} + e_{257}^{*}$$

$$\xrightarrow{h} 1 + 4^{-1}e_{123}^{*} + (4\lambda)^{-1}(e_{156}^{*} - e_{246}^{*} + e_{345}^{*}).$$

When $p \neq \pm 2$, apply the same action which we have used in (a) and (b) followed by the action of $h \in H$ such that

$$\begin{split} \phi(h)e_1 &= -2\lambda(1-\omega^2)^{-1}e_4\,, \qquad \phi(h)e_2 = 2(1-\omega^2)e_2\,, \\ \phi(h)e_3 &= \left(\sqrt{p^2-4}\right)^{-1}e_3\,, \quad \phi(h)e_4 = e_1\,, \quad \phi(h)e_5 = e_5\,, \\ \phi(h)e_6 &= (4\lambda)^{-1}e_6\,, \qquad \phi(h)e_7 = \left(\sqrt{p^2-4}\right)e_7. \end{split}$$

We obtain

$$1 + \frac{(1 - \omega^2)^{-1}}{2} e_2 e_3 e_4 e_7 + \frac{(1 - \omega^2)}{2\lambda} e_1 e_3 e_5 e_7 + (\sqrt{p^2 - 4}) e_7^*$$

$$\xrightarrow{h} 1 + e_1 e_2 e_3 e_7 + e_3 e_4 e_5 e_7 + e_7^*.$$

Futhermore, if $\sqrt{p^2-4} \in k^*$, with the action of $h \in H$ such that $\phi(h)e_i = e_i$ for $i \neq 3$, 7, and $\phi(h)e_3 = (\sqrt{p^2-4})^{-1}e_3$, and $\phi(h)e_7 = (\sqrt{p^2-4})e_7$, then spinors (12) and (14) are equivalent to spinors (1) and (6), respectively.

If $\sqrt{p^2-4} \notin k^*$, spinor (12) is equivalent to

$$1 + e_{147}^* + e_{257}^* - \frac{\lambda}{4} e_{367}^* \tag{16}$$

as shown above and spinor (14) is equivalent to spinor (15) by the action of

$$h_4h_3s_{5,6}(-1)h_2s_{11,14}(-1)s_{12,13}(-1)s_{4,7}(-1)h_1$$
,

where $h_i \in H$, i = 1, 2, 3, 4, such that

$$\phi(h_i) = \begin{bmatrix} \alpha_i & 0 \\ 0 & {}^t\alpha^{-1} \end{bmatrix},$$

with

Consolidating these results, we have

Proposition. Every spinor of the form $1 + x_4 + x_6$ with $x_6 \neq 0$ is equivalent to one of the following spinors:

(1)
$$1 + e_7^*$$
,

(16)
$$1 + e_{147}^* + e_{257}^* - \frac{\lambda}{4} e_{367}^*,$$

(6)
$$1 + e_1 e_2 e_3 e_7 + e_7^*,$$

(15)
$$1 + 4^{-1}e_{123}^* + (4\lambda)^{-1}(e_{156}^* - e_{246}^* + e_{345}^*),$$

(7)
$$p(1 + e_1e_2e_3e_7 + e_4e_5e_6e_7 + e_7^*),$$

(8)
$$p(1+4^{-1}e_{123}^*+(4\lambda)^{-1}(e_{156}^*-e_{246}^*+e_{345}^*)-e_7^*),$$
(9)
$$1+e_1e_2e_3e_7+e_3e_4e_5e_7+e_7^*,$$

$$(9) 1 + e_1 e_2 e_3 e_7 + e_3 e_4 e_5 e_7 + e_7^*$$

$$(10) 1 + e_1e_2e_3e_7 + e_4e_3e_5e_7 + e_6e_5e_2e_7 + e_7^*,$$

$$(11) 1 + e_{147}^* + e_{257}^*,$$

$$(13) 1 + e_1 e_2 e_3 e_7 + e_{147}^* + e_{257}^*,$$

where $p \in k^*$ and $\lambda \in \mathbf{F}_q^* - (\mathbf{F}_q^*)^2$, with q a power of any odd prime number.

Now we consider the case when $x_6 = 0$. Let

$$\sigma \colon \Lambda^3(\text{Aff}^7) \to \Lambda^4(\text{Aff}^7)$$

be the k-linear mapping with respect to the basis e_i, e_i, e_i , $1 \le i_1 < i_2 < i_3 \le i_1 < i_2 < i_3 \le i_1 < i_2 < i_3 \le i_3 < i_3 <$ 7, such that $\sigma(e_{i_1}e_{i_2}e_{i_3})=e_{i_1i_2i_3}^*$. Since the action of SL_7 on $\Lambda^3({\rm Aff}^7)$ and $\Lambda^4(Aff^7)$ are contragredient and since

$$\Lambda^{3}(Aff^{7}) \xrightarrow{h \in SL_{7}} \Lambda^{3}(Aff^{7})$$

$$\downarrow \sigma \qquad \qquad \downarrow \sigma$$

$$\Lambda^{4}(Aff^{7}) \xrightarrow{{}^{\iota}h^{-1} \in SL_{7}} \Lambda^{4}(Aff^{7})$$

commutes, any two elements $u, v \in \Lambda^3(Aff^7)$ are SL_7 -equivalent if and only if $\sigma(u), \sigma(v) \in \Lambda^4(Aff^7)$ are SL_7 -equivalent. Using Igusa's [11] classification of the trivectors of seven-dimensional space over finite fields, every element of $\Lambda^3(Aff^7)$ is SL_7 -equivalent to one of the following trivectors:

0;
$$e_5e_6e_7$$
; $e_1e_3e_7 + e_2e_4e_7$;
 $e_1e_4e_7 + e_2e_5e_7 + e_3e_6e_7$;
 $e_2e_3e_4 + e_3e_1e_5 + e_1e_2e_6$;
 $e_1e_2e_3 + e_4e_5e_6$;
 $e_1e_2e_3 + (e_1e_4 + e_2e_5 + e_3e_6)e_7$;
 $e_1e_2e_3 + e_4e_5e_6 + e_1e_4e_7$;
 $e_2e_3e_4 + e_3e_1e_5 + e_1e_2e_6 + (e_1e_4 + e_2e_5 + e_3e_6)e_7$;
 $e_1e_2e_3 + e_4e_5e_6 + (e_1e_4 + e_2e_5 + e_3e_6)e_7$;
 $e_1e_2e_3 + \lambda(e_1e_4e_7 + e_5e_6e_1 + e_6e_4e_2 + e_4e_5e_3)$; or
 $e_1e_2e_3 + \lambda(e_5e_6e_1 + e_6e_4e_2 + e_4e_5e_3)$

with $\lambda \in \mathbb{F}_q^* - (\mathbb{F}_q^*)^2$. We conclude that the spinors of the form $x = 1 + x_4$ are equivalent to one of the following spinors:

```
(17)
(18)
                                1 + e_1 e_2 e_3 e_4:
                                1 + e_{137}^* + e_{247}^*;
(19)
                                1 + e_{147}^* + e_{257}^* + e_{367}^*;
(20)
                                1 + e_{234}^* - e_{135}^* + e_{126}^*;
(21)
(22)
                                1 + e_{123}^* + e_{456}^*;
                                1 + e_{123}^* + e_{147}^* + e_{257}^* + e_{367}^*;
(23)
(24)
                                1 + e_{123}^* + e_{456}^* + e_{147}^*;
                                1 + e_{234}^* - e_{135}^* + e_{126}^* + e_{147}^* + e_{257}^* + e_{367}^*;
(25)
                                1 + e_{123}^* + e_{456}^* + e_{147}^* + e_{257}^* + e_{367}^*;
(26)
                                1 + e_{123}^* + \lambda(e_{147}^* + e_{156}^* - e_{246}^* + e_{345}^*); or
(27)
                                1 + e_{123}^* + \lambda (e_{156}^* - e_{246}^* + e_{345}^*),
(28)
```

where $\lambda \in \mathbf{F}_q^* - (\mathbf{F}_q^*)^2$, with q a power of any odd prime number.

As in [20] we may claim that the spinors (19)-(26) are equivalent to spinors of the Proposition. Futhermore, spinors (27) and (28) are equivalent to spinors of the type $1 + x_4 + x_6$, $x_6 \neq 0$, with the action of $s_{5,6}(1)s_{2,3}(\lambda)s_{11,14}(1)$. Consequently they are also equivalent to spinors of the Proposition.

We present the orbital decomposition:

Theorem 2. Put

$$\begin{split} &\xi_0 = p(1 + e_1 e_2 e_3 e_7 + e_4 e_5 e_6 e_7 + e_7^*)\,, \\ &\xi_0' = p(1 + 4^{-1} e_{123}^* + (4\lambda)^{-1} (e_{156}^* - e_{246}^* + e_{345}^*) - e_7^*)\,, \\ &\xi_1 = 1 + e_1 e_2 e_3 e_7 + e_4 e_3 e_5 e_7 + e_6 e_5 e_2 e_7 + e_7^*\,, \\ &\xi_5 = 1 + e_1 e_2 e_3 e_7 + e_3 e_4 e_5 e_7 + e_7^*\,, \\ &\xi_{10} = 1 + e_1 e_2 e_3 e_7 + e_7^*\,, \end{split}$$

$$\begin{aligned} \xi_{10}' &= 1 + 4^{-1}e_{123}^* + (4\lambda)^{-1}(e_{156}^* - e_{246}^* + e_{345}^*) \,, \\ \xi_{14} &= 1 + e_1e_2e_3e_7 + e_{147}^* + e_{257}^* \,, \\ \xi_{20} &= 1 + e_7^* \,, \\ \xi_{20}' &= 1 + e_{147}^* + e_{257}^* - \frac{\lambda}{4}e_{367}^* \,, \\ \xi_{21} &= 1 + e_{147}^* + e_{257}^* \,, \\ \xi_{29} &= 1 + e_1e_2e_3e_4 \,, \\ \xi_{42} &= 1 \,, \quad and \\ \xi_{64} &= 0 \,, \end{aligned}$$

where $p \in k^*$ and $\lambda \in \mathbf{F}_q^* - (\mathbf{F}_q^*)^2$, with q a power of any odd prime number. Then

$$X(\mathbf{F}_q) = \bigcup_{\xi_i} (G \cdot \xi_i)(\mathbf{F}_q)$$

for all $\xi_i = \xi_0, \xi_1, \dots, \xi_{64}$ and $\xi'_0, \xi'_{10}, \xi'_{20}$.

3. The fixers

In this section we shall investigate the algebraic structure of the fixers in G. The results are summarized as

Theorem 3. There exists a subgroup H_i of $G_{\xi_i}^0$ for $i = 0, 1, 5, 10, 14, 20, 21, 29, 42 and a subgroup <math>H_i'$ of $G_{\xi_i'}^0$ for i = 0, 10, 20 such that

$$\begin{split} H_0 &\cong G_2 \times G_2 \,, \quad H_0' \cong G_2(\mathbb{F}_{q^2}) \,, \quad H_1 \cong (GL_1 \times G_2) \cdot (G_a)^{14} \,, \\ H_5 &\cong (GL_1 \times SL_2 \times_{\mathbb{Z}_2} SP_4) \cdot U^{19} \,, \quad H_{10} \cong (GL_1 \times SL_3 \times SL_3) \cdot (U^{21}) \,, \\ H_{10}' &\cong (GL_1(\mathbb{F}_q) \times SL_3(\mathbb{F}_{q^2}))(U^{21}(\mathbb{F}_q)) \,, \quad H_{14} \cong (GL_1 \times SL_4) \cdot (U^{26}) \,, \\ H_{20} &\cong (GL_1 \times SL_6) \cdot (G_a^{12}) \,, \quad H_{20}' \cong ((GL_1 \times SU_6)(\mathbb{F}_q))(G_a^{12}(\mathbb{F}_q)) \,, \\ H_{21} &\cong (GL_1 \times SP_6 \times_{\mathbb{Z}_2} G_m) \cdot U^{26} \,, \quad H_{29} \cong (GL_1 \times SL_3 \times Spin_7) \cdot (U^{27}) \,, \\ H_{42} &\cong (GL_7) \cdot (G_a^{21}). \end{split}$$

Since each of the above cases has a similar proof, we shall only demonstrate the procedure for the case of i = 0, 10, 20 when the identity of the group G_{ξ_i} has two connected components.

 H_0 is the subgroup of G generated by

$$T_0 = T \cap G_{\xi_0} = \left\{ t = \prod_{i=1}^7 g_i(t_i); \prod_{i=1}^6 t_i = t_1 t_2 t_3 t_7 = 1, \atop t_4 t_5 t_6 t_7 = 1, t_1, \dots, t_7 \in G_m \right\}$$

and by elements:

$$s_{1,9}(\lambda)$$
, $s_{2,8}(\lambda)$, $s_{1,10}(\lambda)$, $s_{3,8}(\lambda)$, $s_{2,10}(\lambda)$, $s_{3,9}(\lambda)$, $s_{4,12}(\lambda)$, $s_{5,11}(\lambda)$, $s_{4,13}(\lambda)$, $s_{6,11}(\lambda)$, $s_{5,13}(\lambda)$, $s_{6,12}(\lambda)$, $s_{1,2}(\lambda)s_{7,10}(\lambda)s_{10,14}(\lambda)$, $s_{8,9}(\lambda)s_{3,7}(\lambda)s_{3,14}(-\lambda)$, $s_{2,7}(\lambda)s_{8,10}(-\lambda)s_{2,14}(-\lambda)$, $s_{9,14}(-\lambda)s_{1,3}(\lambda)s_{7,9}(-\lambda)$, $s_{2,3}(\lambda)s_{7,8}(\lambda)s_{8,14}(\lambda)$, $s_{9,10}(\lambda)s_{1,14}(-\lambda)s_{1,7}(\lambda)$, $s_{6,7}(-\lambda)s_{6,14}(-\lambda)s_{11,12}(-\lambda)$, $s_{13,14}(-\lambda)s_{7,13}(\lambda)s_{4,5}(-\lambda)$, $s_{5,7}(-\lambda)s_{5,14}(-\lambda)s_{11,13}(\lambda)$, $s_{12,14}(\lambda)s_{7,12}(-\lambda)s_{4,6}(-\lambda)$, $s_{5,6}(-\lambda)s_{7,11}(\lambda)s_{11,14}(-\lambda)$, and $s_{12,13}(-\lambda)s_{4,14}(-\lambda)s_{4,7}(-\lambda)$,

where $\lambda \in G_a$. One can easily verify that $H_0 \subset G_{\xi_0}^0$ by examining directly each of the elements listed above. Futhermore, by the lemma in §1, we have

$$\phi(H_0) = \left\{ \begin{bmatrix} \alpha & \beta \\ \gamma & \delta \end{bmatrix} \right\} \cong \left\{ (A) \cdot (B) \right\} \cong G_2 \cdot G_2,$$

where

$$\alpha = \begin{bmatrix} X & 0 & b \\ & & c \\ & & d \\ 0 & Y & e \\ & & f \end{bmatrix}, \quad X = (x_{ij}), Y = (y_{ij}) \in SL_3,$$

$$\beta = \begin{bmatrix} 0 & t & -s & & & -a \\ -t & 0 & r & & 0 & & -b \\ s & -r & 0 & & & & -c \\ & & 0 & w & -v & d \\ 0 & & -w & 0 & u & e \\ & & v & -u & 0 & f \\ a & b & c & -d & -e & -f & 0 \end{bmatrix},$$

$$\gamma = \begin{bmatrix} 0 & -c & b & & & & r \\ c & 0 & -a & & 0 & & s \\ -b & a & 0 & & & t \\ & & 0 & f & -e & -u \\ 0 & & -f & 0 & d & -v \\ & & e & -d & 0 & -w \\ -r & -s & -t & u & v & w & 0 \end{bmatrix},$$

$$\delta = \begin{bmatrix} x_{11}^{-1} & -x_{21} & -x_{31} & 0 & 0 & 0 & -r \\ -x_{12} & x_{22}^{-1} & -x_{32} & 0 & 0 & 0 & -s \\ -x_{13} & -x_{23} & x_{33}^{-1} & 0 & 0 & 0 & -t \\ 0 & 0 & 0 & y_{11}^{-1} & -y_{21} & -y_{31} & -u \\ 0 & 0 & 0 & -y_{12} & y_{22}^{-1} & -y_{32} & -v \\ 0 & 0 & 0 & -y_{13} & -y_{23} & y_{31}^{-3} & -w \\ -a & -b & -c & -d & -e & -f & 1 \end{bmatrix},$$

$$A = \begin{bmatrix} 0 & 2r & 2s & 2t & 2a & 2b & 2c \\ a & & & 0 & t & -s \\ b & X & & -t & 0 & r \\ c & & & s & -r & 0 \\ r & 0 & -c & b & & & \\ s & c & 0 & -a & & -{}^{t}X & \\ t & -b & a & 0 & & & \end{bmatrix},$$

$$B = \begin{bmatrix} 0 & 2u & 2v & 2w & 2d & 2e & 2f \\ d & & & 0 & w & -v \\ e & Y & & -w & 0 & u \\ f & & v & -u & 0 \\ u & 0 & f & -e \\ v & -f & 0 & d & & -{}^{t}Y \\ w & e & -d & 0 \end{bmatrix}.$$

Since G_2 has no center, $H_0 \cong G_2 \times G_2$.

Similarly we can construct the elements that generate H_{10} and H_{20} and show that $G^0_{\xi_{10}}\supset H_{10}\cong (GL_1\times SL_3\times SL_3)\cdot (U^{21})$ and $G^0_{\xi_{20}}\supset H_{20}\cong (GL_1\times SL_6)\cdot (G^{12}_a)$. Now let $g_0=g_{10}h\in G(\mathbf{F}_{q^2})$ with $g_{10}\in H$ (see §2) such that

$$\phi(g_{10}) = \begin{bmatrix} \alpha & 0 \\ 0 & {}^{t}\alpha^{-1} \end{bmatrix}, \qquad \alpha = \begin{bmatrix} I_3 & I_3 & 0 \\ \mu I_3 & -\mu I_3 & 0 \\ 0 & 0 & -(2\mu)^{-3} \end{bmatrix},$$

i.e.,

$$\begin{split} g_{10} = & \{ 1 + \mu e_4 e_{7+1} + (e_1 - (\mu + 1)e_4)e_{7+4} - (\mu + 1)e_1 e_4 e_{7+1}e_{7+4} \} \\ & \cdot \{ 1 + \mu e_5 e_{7+2} + (e_2 - (\mu + 1)e_5)e_{7+5} - (\mu + 1)e_2 e_5 e_{7+2}e_{7+5} \} \\ & \cdot \{ 1 + \mu e_6 e_{7+3} + (e_3 - (\mu + 1)e_6)e_{7+6} - (\mu + 1)e_3 e_6 e_{7+3}e_{7+6} \} \\ & \cdot g_7 (-(2\mu)^{-3}) \,, \end{split}$$

where $\,\mu=\sqrt{\lambda}\,,\,\,\lambda\in {\mathbf F}_q^*-({\mathbf F}_q^*)^2\,,$ and $\,h\in H\,$ such that

$$\phi(h) = \begin{bmatrix} \alpha & 0 \\ 0 & {}^{t}\alpha^{-1} \end{bmatrix}, \qquad \alpha = \begin{bmatrix} I_5 & 0 \\ 0 & -I_2 \end{bmatrix},$$

i.e.,

$$h = g_1(1)g_2(1)g_3(1)g_4(1)g_5(1)g_6(-1)g_7(-1).$$

Let $g_{20} \in G(\mathbf{F}_{q^2})$ such that

$$g_{20} = s_{1,4} \left(-\frac{\sqrt{\mu^2 + 4}}{2} \right) s_{2,5} \left(-\frac{\sqrt{\mu^2 + 4}}{2} \right) s_{10,13} \left(-\frac{\sqrt{\mu^2 + 4}}{2} \right)$$

$$\cdot s_{8,11} (-\omega) s_{9,12} (-\omega) s_1 ((1 - \omega^2)^{-1}) s_{3,6} (\mu \omega^2 (1 - \omega^2)^{-1})$$

$$\cdot s_{1,4} (\omega (1 - \omega^2)) s_{2,5} (\omega (1 - \omega^2)^{-1}) s_{10,13} (-\mu) h',$$

where

$$\omega = \frac{\sqrt{\mu^2 + 4} - \mu}{2}$$

with $\mu = \sqrt{\lambda}$, $\lambda \in \mathbf{F}_q^* - (\mathbf{F}_q^*)^2$, and $h' \in H$ such that

$$\phi(h') = \begin{bmatrix} \alpha & 0 \\ 0 & {}^{t}\alpha^{-1} \end{bmatrix}, \qquad \alpha = \begin{bmatrix} I_5 & 0 & 0 \\ 0 & \mu & 0 \\ 0 & 0 & \mu^{-1} \end{bmatrix},$$

i.e.,

$$h' = g_1(1)g_2(1)g_3(1)g_4(1)g_5(1)g_6(\mu)g_7(\mu^{-1}).$$

Then the action of g_i transforms ξ_i to ξ_i' , i.e., $\rho(g_i) \cdot \xi_i = \xi_i'$ for i = 0, 10, 20. Hence for any $h_i \in H_i \subset G_{\xi_i}$ we have

$$\rho(g_{i}h_{i}g_{i}^{-1}) \cdot \xi_{i}' = \rho(g_{i}h_{i}g_{i}^{-1}) \cdot \rho(g_{i})\xi_{i} = \rho(g_{i}h_{i}) \cdot \rho(g_{i})^{-1}\rho(g_{i})\xi_{i}$$
$$= \rho(g_{i}h_{i}) \cdot \xi_{i} = \rho(g_{i})(\rho(h_{i}) \cdot \xi_{i}) = \rho(g_{i}) \cdot \xi_{i} = \xi_{i}'$$

for i=0, 10, 20. Hence $H_i'=g_iH_ig_i^{-1}\subset G_{\xi_i'}^0$ and

$$H'_{0} \cong \phi(H'_{0}) = \phi(g_{0})\phi(H_{0})\phi(g_{0})^{-1} \cong G_{2}(\mathbf{F}_{q^{2}}),$$

$$H'_{10} \cong \phi(H'_{10}) = \phi(g_{10})\phi(H_{10})\phi(g_{10})^{-1} \cong (GL_{1}(\mathbf{F}_{q}) \times SL_{3}(\mathbf{F}_{q^{2}}))(U^{21}(\mathbf{F}_{q})),$$

$$H'_{20} \cong \phi(H'_{20}) = \phi(g_{20})\phi(H_{20})\phi(g_{20})^{-1} \cong ((GL_{1} \times SU_{6})(\mathbf{F}_{q}))(G_{a}^{12}(\mathbf{F}_{q})).$$

4. The classification

First we shall list some formulas for calculating the cardinalities of certain algebraic groups:

$$\begin{split} |SL_n(\mathbf{F}_q)| &= q^{n^2-1} \prod_{1 < i \le n} (1-q^{-i}), \\ |SU_n(\mathbf{F}_q)| &= q^{n^2-1} \prod_{1 < i \le n} (1-(-q)^{-i}), \\ |SO_{2n+1}(\mathbf{F}_q)| &= |SP_{2n}(\mathbf{F}_q)| = q^{n(2n+1)} \prod_{1 \le i \le n} (1-q^{-2i}), \\ |SO_{2n}(\mathbf{F}_q)| &= q^{n(2n-1)} \prod_{1 \le i \le n-1} (1-q^{-2i}) \cdot (1-q^{-n}), \\ |G_2(\mathbf{F}_q)| &= q^{14} (1-q^{-2}) (1-q^{-6}), \\ |U^n(\mathbf{F}_q)| &= |G_a^n(\mathbf{F}_q)| = q^n, \\ |GL_1(\mathbf{F}_q)| &= |G_m(\mathbf{F}_q)| = q (1-q^{-1}), \\ |\mathrm{Spin}_n(\mathbf{F}_q)| &= |SO_n(\mathbf{F}_q)|. \end{split}$$

By Theorem 2 in $\S 2$, we have

$$|X(\mathbf{F}_q)| \leq \sum_{\xi_i} |G \cdot \xi_i(\mathbf{F}_q)|$$

for $\xi_i = \xi_0, \, \xi_1, \, \dots, \, \xi_{64}$ and $\xi_0', \, \xi_{10}', \, \xi_{20}'$. Futhermore,

$$|G \cdot \xi_i(\mathbf{F}_q)| = \frac{|G(\mathbf{F}_q)|}{|G_{\mathcal{E}_i}(\mathbf{F}_q)|}$$

for $\xi_i = \xi_1, \xi_5, \xi_{14}, \xi_{21}, \xi_{29}, \xi_{42}, \xi_{64}$, and

$$|G \cdot \xi_i(\mathbf{F}_q)| = \frac{|G(\mathbf{F}_q)|}{2|G_{\xi_i}(\mathbf{F}_q)|}$$

for $\xi_i = \xi_0$, ξ_0' , ξ_{10} , ξ_{10}' , ξ_{20} , ξ_{20}' . By Theorem 3 in §3, we have $|G_{\xi_i}(\mathbf{F}_q)| \ge |H_i(\mathbf{F}_q)|$ for all ξ_i and ξ_i' . Put $(r) = 1 - q^{-r}$, and $(r)_+ = 1 + q^{-r}$. Then

$$\begin{split} |G \cdot \xi_0(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{2|H_0(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{2|(G_2 \times G_2)(\mathbf{F}_q)|} \\ &= \frac{1}{2}q^{64}(1)(2)_+(6)_+(7)(8)(10) \,, \\ |G \cdot \xi_0'(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{2|H_0'(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{2|G_2(\mathbf{F}_{q^2})|} \\ &= \frac{1}{2}q^{64}(1)(2)(6)(7)(8)(10) \,, \\ |G \cdot \xi_1(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{|H_1(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{|(GL_1 \times G_2)(\mathbf{F}_q) \cdot G_a^{14}(\mathbf{F}_q)|} \\ &= q^{63}(4)(7)(8)(10)(12) \,, \\ |G \cdot \xi_5(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{|H_5(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{|(GL_1 \times SL_2 \times \mathbf{z}_2 SP_4)(\mathbf{F}_q) \cdot U^{19}(\mathbf{F}_q)|} \\ &= q^{59}(2)_+(4)_+(6)(7)(10)(12) \,, \\ |G \cdot \xi_{10}(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{2|H_{10}(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{2|(GL_1 \times SL_3 \times SL_3)(\mathbf{F}_q) \cdot U^{21}(\mathbf{F}_q)|} \\ &= \frac{1}{2}q^{54}(2)_+(3)_+^2(6)_+(7)(8)(10) \,, \\ |G \cdot \xi_{10}'(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{2|H_{10}'(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{2|(GL_1(\mathbf{F}_q) \times SL_3(\mathbf{F}_{q^2})) \cdot U^{21}(\mathbf{F}_q)|} \\ &= \frac{1}{2}q^{54}(2)(7)(8)(10)(12) \,, \end{split}$$

$$\begin{split} |G \cdot \xi_{14}(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{|H_{14}(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{|GL_4(\mathbf{F}_q) \cdot U^{26}(\mathbf{F}_q)|} \\ &= q^{50}(3)_+(7)(8)(10)(12)\,, \\ |G \cdot \xi_{20}(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{2|H_{20}(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{2|(GL_1 \times SL_6)(\mathbf{F}_q) \cdot G_a^{12}(\mathbf{F}_q)|} \\ &= \frac{1}{2}q^{44}(3)_+(5)_+(6)_+(7)(8)\,, \\ |G \cdot \xi_{20}'(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{|H_{20}'(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{|(GL_1 \times SU_6)(\mathbf{F}_q) \cdot G_a^{12}(\mathbf{F}_q)|} \\ &= \frac{1}{2}q^{44}(3)(5)(6)_+(7)(8)\,, \\ |G \cdot \xi_{21}(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{|H_{21}(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{|(GL_1 \times SP_6 \times_{\mathbf{Z}_2} G_m)(\mathbf{F}_q) \cdot U^{26}(\mathbf{F}_q)|} \\ &= q^{43}(1)_+(2)_+(4)_+(7)(10)(12)\,, \\ |G \cdot \xi_{29}(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{|H_{29}(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{|(GL_1 \times SL_3 \times \mathrm{Spin}_7)(\mathbf{F}_q) \cdot U^{27}(\mathbf{F}_q)|} \\ &= q^{35}(2)_+(3)_+(4)_+(6)_+(7)(10)\,, \\ |G \cdot \xi_{42}(\mathbf{F}_q)| &\leq \frac{|G(\mathbf{F}_q)|}{|H_{42}(\mathbf{F}_q)|} = \frac{|G\mathrm{Spin}_{14}(\mathbf{F}_q)|}{|GL_7(\mathbf{F}_q) \cdot G_a^{21}(\mathbf{F}_q)|} = q^{22}(3)_+(5)_+(6)_+(8)\,, \end{split}$$

Adding these together, we get

 $|G \cdot \xi_{64}(\mathbf{F}_q)| = \frac{|G(\mathbf{F}_q)|}{|G_{F_n}(\mathbf{F}_q)|} = \frac{|G\operatorname{Spin}_{14}(\mathbf{F}_q)|}{|G\operatorname{Spin}_{14}(\mathbf{F}_q)|} = 1.$

$$\sum_{\xi_i} |G \cdot \xi_i(\mathbf{F}_q)| \le q^{64} = |X(\mathbf{F}_q)|.$$

Hence we verified that

$$\sum_{\xi_i} |G \cdot \xi_i(\mathbf{F}_q)| = q^{64} = |X(\mathbf{F}_q)|.$$

Since the G-orbits $G \cdot \xi_i$ for all ξ_i are distinct and every G-orbit in X has an \mathbf{F}_q -rational point at least for a high power q of any prime number, we conclude that there are no other G-orbits in X, hence

$$X = \coprod_{\xi_i} G \cdot \xi_i,$$

and $G_{\xi_i} = H_i$ for all ξ_i . We summarize with

Theorem 4. If q is a power of any odd prime, then

$$X(\mathbf{F}_q) = \coprod_i G \cdot \xi_i(\mathbf{F}_q), \quad \operatorname{codim}(G \cdot \xi_i) = i,$$

for all ξ_i in Table 1, where $G_{\xi_i}(\mathbf{F}_q)$ denotes the fixer of spinor ξ_i with respect to $G(\mathbf{F}_q)$ and $G_{\xi_i}^0(\mathbf{F}_q)$ denotes the connected component of the identity of the group $G_{\xi_i}(\mathbf{F}_q)$.

TABLE

į	ζ_i	$G^0_{\xi_i}(\mathbf{F}_q)$	$[G_{oldsymbol{arxeta}_i}(\mathbf{F}_q):G_{oldsymbol{arxeta}_i}^0(\mathbf{F}_q)]$
0	$\mu(1+e_1e_2e_3e_7+e_4e_5e_6e_7+e_7^*),\mu\in {f F}_q^*$	$(G_2 \times G_2)(\mathbf{F}_q)$	2
0	$\mu(1+4^{-1}e_{123}^*+(4\lambda)^{-1}(e_{156}^*-e_{246}^*+e_{345}^*)-e_7^*),$	$G_2(\mathbf{F}_{q^2})$	2
	$\mu \in \mathbb{F}_q^* \ , \ \lambda \in \mathbb{F}_q^* - (\mathbb{F}_q^*)^2$		
_	$1 + e_1e_2e_3e_7 + e_4e_3e_5e_7 + e_6e_5e_2e_7 + e_7^*$	$(G\!L_1\times G_2)(\mathbf{F}_q)\cdot G_a^{14}(\mathbf{F}_q)$	1
5	$1 + e_1e_2e_3e_7 + e_3e_4e_5e_7 + e_7^*$	$(GL_1 \times SL_2 \times_{\mathbf{Z}_2} SP_4)(\mathbf{F}_q) \cdot U^{19}(\mathbf{F}_q)$	1
10	$1 + e_1e_2e_3e_7 + e_7^*$	$(GL_1 \times SL_3 \times SL_3)(\mathbf{F}_q) \cdot U^{21}(\mathbf{F}_q)$	2
10	$1 + 4^{-1}e_{123}^* + (4\lambda)^{-1}(e_{156}^* - e_{246}^* + e_{345}^*),$	$(GL_1(\mathbf{F}_q) \times SL_3(\mathbf{F}_{q^2})) \cdot U^{21}(\mathbf{F}_q)$	2
	$\lambda \in \mathbf{F}_q^* - (\mathbf{F}_q^*)^2$		
14	$1 + e_1e_2e_3e_7 + e_{147}^* + e_{257}^*$	$GL_4(\mathbf{F}_q)\cdot U^{26}(\mathbf{F}_q)$	1
20	$1 + e_7^*$	$(GL_1 \times SL_6)(\mathbf{F}_q) \cdot G_a^{12}(\mathbf{F}_q)$	2
20	$1 + e_{147}^* + e_{257}^* - 4^{-1}\lambda e_{367}^*, \lambda \in \mathbf{F}_q^* - (\mathbf{F}_q^*)^2$	$(GL_1 \times SU_6)(\mathbf{F}_q) \cdot G_a^{12}(\mathbf{F}_q)$	2
21	$1 + e_{147}^* + e_{257}^*$	$(GL_1 \times SP_6 \times_{\mathbf{Z}_2} G_m)(\mathbf{F}_q) \cdot U^{26}(\mathbf{F}_q)$	1
29	$1 + e_1e_2e_3e_4$	$(GL_1 \times SL_3 \times \operatorname{Spin}_7)(\mathbf{F}_q) \cdot U^{27}(\mathbf{F}_q)$	1
42	1	$GL_7(\mathbf{F}_q) \cdot G_a^{21}(\mathbf{F}_q)$	1
64	0	$GSpin_{14}(\mathbf{F}_q)$	1

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